Understanding differential patterns in coral reef recovery: chronic hydrodynamic disturbance as a limiting mechanism for coral colonization

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ABSTRACT: Coral reefs are subject to numerous physical disturbances, and post-disturbance coral recovery potential depends on subsequent re-colonization of impacted habitat. We examined divergent recovery trajectories at 2 proximal reefs disturbed by ship groundings that resulted in clearly delineated areas of altered substrate. Post-disturbance measurements of coral recruitment, survival, and changes in community structure were made approximately annually from 2009–2013 in undisturbed reference areas as well as disturbed pavement and rubble areas. Despite similar initial physical disturbances, there were marked differences between sites, with higher coral recruitment and survival on disturbed pavement than rubble bottom, reference reef, or restoration structures. Subsequent episodic disturbances from rubble mobilization could be a mechanism driving the divergent recovery patterns. To estimate whether local hydrodynamic conditions were sufficient to mobilize rubble, we used a combination of long-term monitoring, hydrodynamic modeling, and rubble transport mechanics to hindcast the potential for substrate mobility. Long-term model simulations of hydrodynamic forcing at the study sites showed multiple events where bottom-orbital velocities exceeded thresholds required to mobilize rubble via sliding or overturning. Our analyses indicate that wave energy mobilizes rubble substrate multiple times annually and suggests a physical limitation on survival of coral recruits relative to those on pavement substrate. Continued mobilizations may establish a positive feedback loop in which continued rubble clast mobilizations cause additional mechanical erosion or breakage and a shift to smaller rubble sizes that would subsequently mobilize at a lower level of hydrodynamic forcing and thus become subject to more frequent and sustained disturbances. The combination of multiple hydrodynamic disturbances and unstable substrate limits coral recovery and thus contributes to prolonged habitat loss.

KEY WORDS: Coral recruitment · Coral reefs · Disturbance recovery · Rubble mobilization · Waves · Ship groundings · Puerto Rico

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Publisher: Inter-Research · www.int-res.com
INTRODUCTION

Habitat degradation is a major contributor to the decline of coral reefs worldwide (Hughes et al. 2003, Hoegh-Guldberg et al. 2007). Coral reef structural complexity can be degraded by both human activities and climate-related stressors. Anthropogenic disturbances that can obliterate reef structure are often related to proximal human uses and include ship groundings (Jaap 2000, Riegl 2001) and blast fishing (Edinger et al. 1998). Loss of reef structure can also be caused by climate change related ecological disturbances, such as widespread coral death from severe bleaching events (Sheppard et al. 2002, 2005, Graham et al. 2006). Future climate change scenarios (e.g. ocean acidification) have been projected to degenerate reefs into rubble banks (Hoegh-Guldberg et al. 2007). Acute, episodic disturbances such as wave forcing from tropical cyclones (Stoddart 1969, Woodley et al. 1981, Harmelin-Vivien 1994) and tsunamis (Scheffers et al. 2009) can reduce the structural complexity of large sections of reef into extensive rubble fields. Repeated or multiple disturbances can also compound effects and slow recovery (e.g. Rasser & Riegl 2002, Wakeford et al. 2008). The physical legacy of disturbance (e.g. degradation of reef into loose rubble) may result in prolonged habitat loss if colonization of new reef-building scleractinian species is unsuccessful.

Scleractinian corals thrive under a range of hydrodynamic conditions, although success varies by species, morphology, and geography. As sessile filter feeders, flow conditions affect prey capture (Sebens & Johnson 1991, Piniak 2002), photosynthesis and respiration (Dennison & Barnes 1988), and removal of waste products (Mass et al. 2010). Flow variability can also mediate corallivory, increasing net growth of corals (Lenihan et al. 2015). Exceeding normal hydrodynamic conditions, water motion can act as a disturbance to benthic marine organisms such as corals (Dayton 1971, Denny 1995). Wave forcing can cause mechanical breakage of colonies or dislodgment from substrate (Woodley et al. 1981, Connell et al. 1997, Madin 2005), and can limit the spatial distributions of coral species on reefs (Storlazzi et al. 2005). Effects of large disturbances are well known to have long-lasting impacts on coral communities (Done 1992, Connell 1997, Connell et al. 1997). After a disturbance, successful recruitment and survival of scleractinian corals are key limitations to the potential for future coral community development (Lirman & Miller 2003) and reef complexity (Steneck et al. 2009); as foundation species, scleractinian corals create habitat through development of biogenic skeletal structure (Dayton 1972, Ellison et al. 2005, Bertness et al. 2006). Many factors can limit the potential for successful coral recruitment (Ritson-Williams et al. 2009), particularly the availability of suitable settlement substrate (e.g. consolidated hard-bottom, Pearson 1981). Although it is known that intervals between extreme events regulate coral community trajectories and return or reset communities to alternate or early successional stages (Done 1999, Hughes & Connell 1999), less is known about how interactions between chronic hydrodynamic stress and reef structure (or lack of structure) may constrain the potential of a disturbed reef to recover to a coral-dominated state.

We used a comparative and mechanistic approach to examine coral colonization and recovery trajectories within areas of localized reef destruction after ship grounding impacts. Coral recruitment and survival were compared between pavement and rubble disturbance sites relative to artificial structures and reference areas for 6 yr. Differential recovery trajectories observed at these sites suggest that physical mechanisms may limit post-disturbance colonization and contribute to observed patterns of colonization success. To explore this, we hindcasted potential limitations on colonization due to the frequency and intensity of hydrodynamic forces, specifically, whether ocean surface wave forcing was sufficient to intermittently mobilize rubble exposed by the initial ship grounding. The specific questions addressed are: (1) After an initial physical disturbance, how does coral colonization differ between exposed substrate, exposed unstable substrate, artificial structures with coral transplants, and a reference reef? (2) How is colonization success limited by abiotic factors, including near-bottom wave-orbital velocity and rubble substrate mobilization? (3) What is the threshold bottom-orbital velocity required to mobilize rubble substrate, and how frequently are these hydrodynamic conditions exceeded in the study area? (4) How does the time scale between potential hydrodynamic disturbance events compare with that of coral colonization?

MATERIALS AND METHODS

Study sites

The study was conducted on the southern coastal shelf of Puerto Rico, south of Bahia de Tallaboa (Fig. 1A). The area is primarily exposed to easterly
trade winds and waves from the southeast (Fig. 1B,C). Water depths in the study area are 9–14 m but increase to more than 2000 m approximately 10 km offshore (Fig. 1D). Tanker vessels frequently transit the study area to access a liquefied natural gas terminal in Guayanilla Bay. Between November 2005 and April 2006, 2 vessels grounded on adjacent reefs, crushing and flattening the reef structural complexity within areas of impact. One vessel (TV ‘Sperchios’) disturbed approximately 1970 m² of coral reef habitat into flattened pavement cleared of corals (hereafter termed ‘pavement site’; Fig. 1). The second vessel (TV ‘Margara’) fractured approximately 6910 m² of coral reef habitat, removing corals and the consolidated reef surface to expose geologic unconsolidated clastic rubble primarily composed of fossilized Acropora cervicornis (hereafter termed ‘rubble site’; Fig. 1). Both grounding impact sites had similar depths (10–13 m). A reef site adjacent to the rubble site was used as a reference reef of unimpacted habitat. To test if coral recruitment occurred on available substrate in the geographic area of the study, restoration structures were included in the experimental design. Artificial structures were constructed and located within a subsection of the rubble area in 2006 to stabilize dislodged corals and partially replace structural complexity lost in the grounding impact. These structures (hereafter termed ‘restoration structures’) were constructed from dislodged pieces of reef limestone and concrete, with diameters ranging between 0.2 and 3.0 m and heights between 0.2 and 1.0 m (Fig. 1). Scleractinian corals and octocorals that survived the disturbance impact were transplanted onto restoration structures in 2006.

Benthic surveys

To evaluate changes in coral recovery after these 2 large and proximal disturbance events, we measured in situ coral recruitment annually at multiple levels of
substrate stability and complexity. These levels were:
(1) disturbed reef area at the rubble site; (2) disturbed reef area at the pavement site; (3) restoration structures within the rubble site; and (4) natural reef adjacent to the rubble site (Fig. 2). Within each level, permanent quadrats (25 cm × 25 cm) were visually surveyed (n = 45 in each of rubble disturbance and reference areas, and n = 30 in each of pavement disturbance and restoration areas) approximately annually from 2008−2013. At each field sampling, new scleractinian and octocoral recruits (≥0.5 cm) within quadrats were visually identified by divers by species, size, and location within the permanent quadrat, and diagrammed onto quadrat maps for reference in future relocations. Existing (non-recruit) corals were identified within quadrats at each survey from 2009−2013. Scleractinian corals were identified to species, and octocorals were identified to genus. Survival of each individual was monitored in subsequent surveys by comparing mapped diagrams of new and existing corals. Survival at each site was calculated as the percentage of corals present at each time interval and present at the subsequent sampling. Photographs of benthic quadrats were collected at each survey and used for analyses of potential rubble movement (below) and change in benthic cover (for details see Section 1 in the Supplement at www.intres.com/articles/suppl/m605p135_supp.pdf). Change in cover for corals and octocorals (%) between 2008 and 2013 was regressed against rubble cover (%) to show the relationship of increases in coral and octocoral cover with rubble coverage at rubble, pavement, restoration, and reference sites.

Coral community composition and size distribution were surveyed visually with belt transects located adjacent to quadrats at rubble, reference, and pavement sites. Within the transect area (10 m × 1 m), scleractinian corals were identified to species, and octocorals were identified to genus. Transect surveys in disturbed areas were conducted approximately annually from 2008 to 2013 (rubble disturbance; n = 8; pavement disturbance; n = 4−6) and in reference areas in 2008 (reference adjacent to rubble disturbance; n = 10) and 2011 (reference adjacent to pavement disturbance; n = 5). No transects were surveyed on restoration structures due to the limited size of the structures. The change in octocoral and scleractinian coral density and richness between the first and last study years (2008 and 2013) was compared between the 2 disturbances using a Mann-Whitney U-test.

Rubble mobilization over time was investigated by comparing the change between survey events in rubble clast locations within individual quadrats. In each photograph of permanently located quadrats taken during benthic surveys, 3−5 individual rubble clasts with uniquely identifiable shapes were identified (n = 249) and then visually searched for in the same quadrat in subsequent years to test whether they could be located in approximately the same or different locations within the quadrat.

Fig. 2. Study locations including (A) rubble and (B) pavement areas injured by ship groundings, (C) artificial restoration structures, and (D) proximal reference site. In (A), the edge of the injury is visible at the top left.
Rubble transport model

We modeled the potential mobility of unconsolidated rubble within the disturbed areas using a mechanistic approach to estimate hydrodynamic forces required to initiate rubble motion. This approach has been well developed in the literature to hindcast hydrodynamic forcing required to mobilize boulders during large wave events (Nandasena et al. 2011). Here, a clast is approximated by an idealized cuboid (defined by a-, b-, c-axes) that is acted upon by gravity \( \left(F_g\right) \), lift \( \left(F_l\right) \), drag \( \left(F_d\right) \), inertial \( \left(F_i\right) \), and friction \( \left(F_f\right) \) forces (Fig. 3), where \( F_f \equiv \mu F_N \) and \( \mu \) is the coefficient of static friction. Hydrodynamic forcing exceeding a critical threshold can initiate rubble clast motion by sliding or over-turning. Sliding occurs when (Fig. 3)

\[
F_f(t) + F_d(t) > \mu \left[F_g - F_l(t)\right] \tag{1}
\]

At the initiation of sliding, the forces on the left and right hand side of Eq. (1) are equal. Forces that are functions of water velocity \( (F_d, F_l, \) and \( F_f) \) vary with time \( (t) \) over a wave period \( (T) \). The quadratic drag and lift forces due to waves are in phase with each other, but they are both 90° out of phase with the inertial force. To estimate when rubble is mobilized, we evaluated when the total time-varying forces are maximum. Following Dean & Dalrymple (1991), but extending their derivation to include the lift force, this occurs when

\[
F_{D,m} + \mu F_{L,m} + \frac{F^2_{L,m}}{4(F_{D,m} + \mu F_{L,m})} - \mu F_g = 0 \tag{2}
\]

where the maximum drag force is \( F_{D,m} = \frac{1}{2} \rho_w C_D A_D U_m^2 \), the maximum lift force is \( F_{L,m} = \frac{1}{2} \rho_w C_L A_D U_m^2 \), and the maximum inertial force is \( F_{I,m} = \rho_w C_M V(U_m) \omega \), where \( \rho_w \) is the density of seawater \( (1023 \text{ kg m}^{-3}) \) and \( \omega = 2\pi/T \) is the wave frequency. The normal force \( (F_N) \) can be expressed as \( F_N = \Delta \rho V_g \), where \( \Delta \rho = \rho_i - \rho_w, \rho_i \) is rubble density, \( g \) is acceleration due to gravity, and \( V \) is the rubble volume. \( C_D, C_L, \) and \( C_M \) are dimensionless coefficients of drag, lift, and added mass, respectively; \( A_D \) is the cross-sectional area of a rubble clast perpendicular to the flow and \( A_D \) is the area of a rubble clast parallel to the flow. The water velocity can be decomposed into contributions from currents and waves \( (U_m = U_c + U_w; \) e.g. Grant & Madsen 1979). Substituting these expressions for forces and velocities into Eq. (2), and noting that the currents do not contribute to the inertial force, one can derive an expression for the threshold velocity at which rubble is mobilized via sliding (see Section 4 of the Supplement for full derivation and solution).

Now, from the free-body diagram, rubble overturning occurs when

\[
F_D I_D + F_l I_l + F_l I_l > F_g I_g \tag{3}
\]

Here, we assume that the moment arms for the drag \( (I_D) \) and inertial \( (I_I) \) forces are the same (half the rubble width \( (l_w) \) and that the moment arms for the lift \( (I_l) \) and gravitational \( (I_g) \) forces are the same (half the rubble height \( (l_h) \) and that the moment arms for the lift \( (I_l) \) and gravitational \( (I_g) \) forces are the same (half the rubble width \( (l_w) \) and that the moment arms for the lift \( (I_l) \) and gravitational \( (I_g) \) forces are the same (half the rubble height \( (l_h) \). A derivation similar to the one for the sliding threshold velocity above can be done to obtain the threshold velocities for the time-varying overturning moments (Section 4 in the Supplement).

Coefficients of drag, lift, inertia, and friction were obtained from previously published literature values. Because empirical coefficients are not well constrained for rubble movement, coefficient values were selected near the center of the range in the literature, and sensitivity analyses were conducted with a range of values. For the general case, we used a \( C_M \) of 1.0 and tested a sensitivity range of ±0.5 from literature values for rubble (McDougal & Sulisz 1990). We used a \( C_D \) of 0.6 with a sensitivity range of ±0.2 as per literature values for cylinders, limpets,
plates, and spheres (Denny 1994). We used a $C_L$ of 0.178 (Nott 2003) with a minimum and maximum range of 0.1 and 0.4 (Cheng & Chiew 1998). We used a $\mu$ of 0.6 and varied by ±0.2 for the sensitivity analyses based on literature values from boulder transport on basalt and sand-covered limestone in laboratory studies (Voropayev et al. 2001, Noormets et al. 2004, Goto et al. 2007). Stability of a rubble clast is a function of rubble properties, including size, shape, and density (Rasser & Riegl 2002). Rubble properties required in the above equations were obtained from statistical values obtained from morphology and density measurements of individual clasts collected from the study site (Section 5, Fig. S3, in the Supplement).

The approach above addresses individual rubble clasts. However, in a rubble field, proximity of multiple clasts likely impacts hydrodynamic forces affecting any given clast. For example, in rubble mobilization in storm conditions, clasts that become loose can knock into neighboring clasts, mobilizing rubble that would not necessarily be mobilized in isolation. Alternatively, friction can be enhanced from interlocking or by reduced area exposed to flow due to sheltering from adjacent objects. We did not incorporate friction and mechanical restraint from interlocking into forcing equations here due to the numerical complexity. We addressed a reduced area exposed to flow due to sheltering from adjacent objects by including the blocking term of Storlazzi et al. (2005) that accounts for reduction in $A_D$. This term has a value from 0–1, corresponding to flow totally blocked to no blockage. We applied a value of 0.5 with recognition that a wide array of conditions is likely to exist in situ, and conducted sensitivity tests for values ranging from 0.3 to 0.9 (Fig. S4 in the Supplement).

### Wave model and hindcasts

We quantified temporal variations in hydrodynamic forces at each study site over a 4 yr period using a spatially resolved ocean wave model hindcast verified with field data. Wave characteristics at the sites were computed using SWAN, a third-generation numerical wave model (Booij et al. 1999, Ris et al. 1999). SWAN solves the spectral wave action equation and accounts for wave propagation, wave generation by local winds, dissipation by bottom friction, depth-limited breaking, and water level changes. The model was forced using wind and wave parameters (1 h intervals) from the Caribbean Coastal Ocean Observing System (CarICOOS) buoy southeast of Ponce, Puerto Rico (17.860° N, 66.52°W; Fig. 1). We employed a nested grid approach to permit adequate wave development and propagation from offshore to the complex inner shelf bottom topography. The outer model domain was 50 km × 16 km with 200 × 200 m grid cell resolution, and spanned the 25 km distance between the offshore buoy and the study area (Fig. 1). The nested inner grid was 16 km × 10 km with 40 × 40 m resolution and contained the study sites. Bathymetry was interpolated from the best available high-resolution NOAA hydrographic surveys (ngdc.noaa.gov).

The accuracy of the wave model was assessed by comparison with field measurements from an inten-

<table>
<thead>
<tr>
<th>Site</th>
<th>Location (Lat, Long)</th>
<th>Depth (m)</th>
<th>Instruments and configuration</th>
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</thead>
<tbody>
<tr>
<td>1 Rubble site</td>
<td>17°57.161’N 66°43.876’W</td>
<td>11.0</td>
<td>(A) 1.2 MHz TRDI ADCP, 0.25 m bins&lt;br&gt;1 Hz sample rate with 10 mode-12 subpings, bursting for 17 min every 0.5 h&lt;br&gt;(B) 6 MHz Nortek Vector ADV&lt;br&gt;16 Hz sample rate, bursting for 17 min every hour&lt;br&gt;(C) SeaBird SBE56 thermistor&lt;br&gt;1 Hz continuous sampling</td>
</tr>
<tr>
<td>2 Pavement site</td>
<td>17°57.587’N 66°46.132’W</td>
<td>9.6</td>
<td>(A) 1.2 MHz TRDI ADCP, 0.25 m bins&lt;br&gt;1 Hz sample rate with 10 mode-12 subpings, bursting for 17 min every 0.5 h&lt;br&gt;(B) 6 MHz Nortek Vector ADV&lt;br&gt;16 Hz sample rate, bursting for 17 min every hour&lt;br&gt;(C) SeaBird SBE56 thermistor&lt;br&gt;1 Hz continuous sampling</td>
</tr>
<tr>
<td>3 Central bay</td>
<td>17°57.701’N 66°45.416’W</td>
<td>10.0</td>
<td>(A) 600 kHz TRDI ADCP, 0.5 m bins&lt;br&gt;2 Hz sample rate in Waves mode, bursting for 17 min every 0.5 h</td>
</tr>
</tbody>
</table>
sive month-long deployment (13 June to 17 July 2014) of oceanographic instruments at the study sites (Table 1, Section 3 in the Supplement). A set of 3 oceanographic moorings was deployed on the seafloor at the rubble and pavement disturbance sites and at a central location within the model domain (Fig. 2, Table 1). Each mooring contained an upward-looking acoustic Doppler current profiler (ADCP) that burst sampled each hour measuring velocity and pressure (see Table 1 for details) to obtain bulk wave parameters of significant wave height ($H_s$), peak wave period ($T_p$), and peak wave direction ($D_p$) following the method of Terray et al. (1990). High-resolution near-bottom velocity measurements were made using acoustic Doppler velocimeters (ADVs) deployed at the rubble and pavement disturbance sites to compute near-bottom wave-orbital velocities ($u_b$). ADV velocity components were despiked following (Islam & Zhu 2013), and $u_b$ values were computed using linear wave theory and integrals of the velocity spectra (Dean & Dalrymple 1991, Wiberg & Sherwood 2008). Numerical model predictions of $H_s$ were compared to in situ observations ($n = 890$ h measurements over the month-long field campaign) for individual sites (Stations 1–3, Fig. 1, Table 1) and for the mean of all sites using linear regression, root mean square error (RMSE), as well as an average skill metric for the 3 (N) sites (Willmott 1982, Hoeke et al. 2011):

$$\text{skill} = \frac{1}{N} \sum_{i=1}^{N} \left[ 1 - \frac{\sum |H_s^{\text{model}} - H_s^{\text{obs}}|^2}{\sum (|H_s^{\text{model}} - H_s^{\text{obs}}|^2 + |H_s^{\text{obs}} - H_s^{\text{obs}}|^2)} \right]$$

where skill varies between 0 and 1, with 1 indicating perfect agreement. Model skill metrics for $u_b$ were computed using the same approach. Modeled $H_s$ and $u_b$ showed good agreement with observations ($H_s$ model skill = 0.90, $u_b$ model skill = 0.83, Fig. 4, Table 2) and thus the model was deemed sufficiently accurate to be used to estimate the potential range of hydrodynamic forcing at the disturbance sites. The wave model was then run for a 4 yr hindcast period (2010–2014) to simulate a representative range of wave conditions at the rubble and pavement disturbance sites.

To address the combined forcing of both waves and currents, we first assumed that the near-bottom flow can be described by a log-layer. We estimated the shear velocity ($u_s$) from short-term field measurements of the depth-averaged velocity (e.g. Buckley et al. 2012)

$$u_s = \frac{1}{\kappa H} \int_{z_0}^{H} \ln \left( \frac{z + z_0}{z_0} \right) dz$$

where $U$ is the depth-averaged velocity, $\kappa$ is the von Karman constant, $H$ is the total water column height, $z$ is the height above the sea floor, and $z_0$ is the roughness height. Assuming that the rubble occupies some layer above the bottom, the layer-averaged velocity ($u_k$) acting on the rubble is

$$u_k = \frac{u_s}{\kappa (z_{k-1} - z_k)} \int_{z_k}^{z_{k-1}} \ln \left( \frac{z + z_0}{z_0} \right) dz$$

Fig. 4. Comparisons between modeled and observed (A) significant wave height ($H_s$) at the 3 study sites and (B) near-bottom-orbital velocity ($u_b$) at the rubble and pavement sites. Model output and observed data are for the intensive in situ oceanographic measurement periods from 13 June to 17 July 2014. Reference lines with an intercept of 0 and slope of 1 are in grey.
Mean depth-averaged speeds derived from short-term ADCP measurements were 0.05 and 0.10 m s\(^{-1}\) for the pavement and rubble sites, respectively. Assuming a \(z_0\) value of 0.01 m gives \(u_\ast\) values of 0.003 and 0.007 m s\(^{-1}\), respectively, so from Eqs. (5) & (6), \(u_k\) for the pavement site was 0.02 and 0.04 m s\(^{-1}\) for the rubble site. These values were included as constant \(U_c\) values in the rubble model forcing equations.

To determine the frequency of occurrence of projected hydrodynamic mobilization events, we compared 4 yr model hindcast estimates of bottom-orbital velocity at the rubble and pavement sites to time-dependent threshold velocities required to mobilize rubble through sliding or overturning (Eqs. 2 & 3) using our measured rubble properties. We identified when potential mobilization forces exceeded threshold bottom-orbital velocities at the rubble injury site and calculated time intervals between mobilization events. Only events more than 4 d apart were included to ensure that the interval estimates represented distinct disturbance events.

### RESULTS

**Coral recruitment, survival, and community change**

Successful colonization of corals within the rubble site was limited by low recruitment, survival, and cover (Figs. 5, 6A,C, & 7B). Overall change in benthic cover of scleractinian and octocorals during the study was related to the amount of rubble cover (Fig. 5). Density and survival of scleractinian coral recruits were consistently lower at the rubble site than on the reference reef or on restoration structures during every year of the study (Figs. 5, 6A,C). As a result, coral density and species richness increased at the pavement site from 2008 to 2013, as recruits survived and grew into the larger size class (Figs. 7A,C, 8), in contrast to the rubble transects where scleractinian recruits remained in the small (<5 cm) size class (Fig. 7B). The change in scleractinian density during the study was also significantly higher within the pavement site than the rubble site (scleractinian rubble and pavement mean ranks were 7.8 and 13.2, respectively; \(U = 0, Z = -2.6, p < 0.05, r = -0.8\); Fig. 7). Scleractinian recruitment was dominated in the rubble disturbance by *Siderastrea siderea*, and in the pavement disturbance by *S. siderea* and *Porites astreoides* (Fig. 9). Scleractinian species richness averaged 9.3 within the rubble site and 10.2 at the pavement site, comparable to the 10 species observed in the restoration site, but fewer than the 22 species observed in the reference site (Fig. 9). The change in species richness from 2008 to 2013 did not differ significantly between pavement and rubble disturbances for either scleractinians or octocorals (Figs. 8 & 9). Therefore, the observed differences in coral density and survival between the pavement...
Fig. 6. Comparison of coral colonization success between rubble, pavement, restoration, and reference sites showing (A,B) density (error bars represent SE) and (C,D) survival of recruits. Recruit survival represents the percentage of corals present at the previous survey.

Fig. 7. Size-frequency histogram for scleractinian corals (A–C) and octocorals (D–F) at rubble (A,D), pavement (B,E), and reference (C,F) sites in 2008 (rubble, pavement, reference) and 2013 (rubble, pavement). Sizes (cm) are size-class midpoints. Error bars represent standard error of the mean.
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and rubble sites are likely not due to changes in community composition over time (Fig. 8), although this was not well resolved in the data due to the limited number of recruits at the rubble site. Patterns in coral density and survival were similar between the reference and restoration areas (Fig. 5), likely due to similarities in benthic cover of turf algae, macroalgae, and other organisms (Section 1, Fig. S1B, in the Supplement).

Octocoral recruit density and survival were higher than patterns seen in scleractinian corals, although still limited and variable. Density and survival of recruits were predominantly lower at the rubble site than the pavement site, reference site, or restoration structures (Figs. 6B, D, & 7). The change in octocoral density from 2008 to 2013 was significantly higher within the pavement site than within the rubble site (rubble and pavement mean ranks were 4.9 and 9.8, respectively; \( U = 0, Z = -2.2, p < 0.05, r = -0.6 \); Fig. 7); this appeared due to increases in density of octocorals sized from 5 to >20 cm at the pavement site (Fig. 6), as surviving octocorals grew larger. In contrast, although the density of octocorals at the rubble site was higher than that of scleractinians (Figs. 7 & 8) and some octocorals grew into larger size classes, octocoral density did not significantly increase from 2008 to 2013 (Fig. 7).

At the rubble site, analysis of benthic quadrat photographs showed that selected individual rubble clasts changed locations within permanent quadrats between field surveys, indicating repeated mobilization. Of 249 rubble clasts tracked between sampling events, only 5 (2.0%) were relocated in the subsequent years’ photos. These 5 consistent pieces appeared larger than the mean of the rubble size distribution, and were in approximately the same location for a single subsequent year.

Fig. 8. Change in (A) density and (B) richness for scleractinians and octocorals from within the rubble and pavement sites. Richness shows change in number of scleractinian species and octocoral genera per 10 m² transect

Fig. 9. Change in species composition from 2008–2013 for impacted coral communities relative to reference populations (Ref) at (A,C) pavement and (B,D) rubble sites. Scleractinian coral composition is in A and B, and octocoral composition is in C and D
Rubble and wave energy models

The 4 yr wave model hindcast showed that wave forcing near the substrate was similar between the rubble and pavement disturbance sites (Fig. 10). Multiple events occurred during the model hindcast period when bottom-orbital velocities exceeded thresholds required to mobilize mean-sized rubble (Figs. 11 & 12). The mean return intervals for forcing required to slide rubble was approximately 7 d (maximum return interval approximately 23 d), shorter than for return intervals of forces required to overturn rubble, where the mean return interval was 12 d (maximum 134 d; Fig. 12A). The duration of sliding events was longer (mean = 1 d, max = 58 d) than for overturning events (mean = 0.5 d, max = 22 d; Fig. 12B). During multiple events each year, forcing was sufficient to slide (60–131 events yr⁻¹) and overturn rubble (151–241 events yr⁻¹). A number of the mobilization events are linked to the passage of tropical storms or hurricanes (Fig. 11). Forcing sufficient to mobilize rubble by sliding or overturning occurred in all months (Figs. 11 & 12C), with peak occurrence in August through December during the Atlantic tropical cyclone season (Fig. 12C).

DISCUSSION

The frequency and spatial extent of multiple disturbances constrain recovery processes (Turner et al. 1993), including colonization of key species (Connell 1997). Temporal intervals between extreme events regulate coral community trajectories and return or reset communities to alternate or early successional stages (Done 1999, Hughes & Connell 1999). Here, we show how diverging coral recovery trajectories are caused by repeated, chronic physical disturbances from wave energy. Coral colonization into 2 proximal ship grounding sites with a similar post-impact recovery time was far higher on a reef flattened to pavement than on reef area broken to rub-

Fig. 10. Spatial distribution of (A,C,E) significant wave heights and (B,D,F) bottom-orbital velocities obtained from long-term wave model hindcasts: (A,B) calmest 5% percentile, (C,D) mean, and (E,F) stormiest 95% conditions.
Fig. 11. Comparison between model hindcast estimates of bottom-orbital velocities ($u_b$) and the critical orbital velocities required to slide ($u_{slide}$) or overturn ($u_{overturn}$) coral rubble at (A) rubble and (B) pavement sites. Grey vertical lines indicate named storm events south of Puerto Rico within 100 km. Solid red and blue lines indicate general values of coefficients used in forcing needed to overturn (red) or slide (blue) rubble. Shading indicates minimum and maximum coefficient sensitivity values for $u_{slide}$ (red) and $u_{overturn}$ (blue).

Fig. 12. (A) Return intervals (in days) between hydrodynamic forcing events that exceeded threshold levels to mobilize coral rubble through overturning (light grey) or sliding (dark grey). Return intervals for threshold conditions to overturn rubble have a longer return interval than those to slide rubble. (B) Durations of hydrodynamic forcing conditions exceeding conditions to mobilize coral rubble through overturning are shorter than those required for sliding. (C) Probability of occurrence of hydrodynamic forcing events by month. Chronic hydrodynamic forcing events of rubble occur in all months and more frequently in the hurricane season. (D) Probability of occurrence of hydrodynamic forcing events by year, showing that threshold conditions were exceeded in all study years.
ble. Recruitment density and survival on pavement substrate also initially exceeded that of the reference and restoration sites, a pattern that is likely indicative of less available benthic space in reference and restoration sites due to existing corals, octocorals, and other benthic organisms. As coral recruits at the pavement site survived and grew into larger size classes, new recruitment declined, also supporting a pattern of space availability limitations as more corals became established. In contrast, within unstable substrate, the data suggest a colonization bottleneck for corals and sustained loss of coral habitat. To contrast these diverging recovery trajectories on reefs with otherwise similar coral species composition, we further explored the potential physical mechanisms behind the observed biological patterns.

Using rubble mechanics, we showed that the hydrodynamic forcing hindcast for the study sites would be sufficient to mobilize rubble on a chronic timescale not limited to large, infrequent events such as named tropical storms. The recruitment limitations within unstable substrate observed in this study are consistent with previous studies that have shown correlations between decreased survival of small corals and substrate mobilization by water motion (Fox et al. 2003, Fox & Caldwell 2006, Yadav et al. 2016). Our mechanistic model indicates that substrate instability lowers a threshold for subsequent hydrodynamic disturbances. Such chronic multiple disturbances limit coral colonization relative to adjacent disturbed areas with comparable hydrodynamic forcing but consolidated substrate. Without stabilization of unconsolidated rubble, the threshold for disturbance remained low and did not recover during the study. In contrast, the threshold for hydrodynamic disturbance at the pavement site was much higher, as evidenced by successful coral colonization. The fate of unstable rubble depositions is determined by frequency and intensity of subsequent hydrodynamic disturbances (Scoffin 1993), and continued mobilizations may establish a positive feedback loop in which continued rubble clast mobilizations cause additional mechanical erosion or breakage and a shift to ever smaller rubble sizes. Smaller rubble clasts subsequently mobilize at a lower level of hydrodynamic forcing, and thus become subject to more frequent and sustained disturbances.

Estimating hydrodynamic forces on submerged objects contains inherent uncertainties. Spatial variability in flow patterns (notably, at a scale smaller than the 40 m model grid cell resolution of this study) may influence probabilities of rubble mobilization. Small-scale flow patterns may drive hydrodynamic patchiness within a disturbance site, resulting in some areas being more likely to mobilize than others. For example, colonies along edges of undamaged reef surrounding disturbance area may have turbulent wakes (Hench & Rosman 2013) that may influence the likelihood of rubble mobilization. Potential for substrate mobility is influenced by friction and drag forces’ sensitivities to rubble exposure to the flow, rubble interlocking, or ratio of rubble to sediment. Improved modeling of rubble mobility could include rubble collisions (Imamura et al. 2008, Nandasena & Tanaka 2013) and rubble interactions with smaller sediment size classes (Kain et al. 2012). Finally, bioturbation (e.g. sand tilefish, rays) may mobilize rubble, but is difficult to quantify on its own or in conjunction with hydrodynamic mobilization.

Alternative explanations for the comparatively lower coral recruitment into the rubble site than the pavement site include differential temperature stresses leading to mortality (e.g. via bleaching), or competition for space with existing benthic organisms such as algae. Mean near-bottom water temperature at the pavement site was slightly warmer than the rubble site, with the same variance, and skewed toward higher temperatures (Section 2, Fig. S2, in the Supplement). One might expect that greater thermal stress at the pavement site could translate to less favorable thermal conditions for coral survival; however, since coral density and survival were higher at the pavement site, it appears that thermal stress is not the dominant physical driver accounting for differences between sites. Cover of other benthic organisms also did not appear to be substantially different between the pavement and rubble sites. Both were similar in cover for benthic algae, turf, and sponges; however, the rubble site had more crustose coralline algae than the pavement site for several years (Figs. S1A & S1B in the Supplement), a condition which would seem favorable for coral settlement (Ritson-Williams et al. 2009). Cementation and encrustation (Perry 1999) and sponge stabilization (Biggs 2013) have been reported to be biological mechanisms to stabilize substrate, and potentially provide suitable substrate for successful multi-species coral colonization (e.g. Dollar & Tribble 1993, Hughes 1999, Perry 2005). In our study, however, only limited stabilization of unconsolidated substrate through biological mechanisms was noted, and biological stabilization was not observed to be a significant contributor to stabilize the rubble site.
Although direct observation of rubble mobilization and coral mortality during hydrodynamic events proved elusive in this study, the findings are consistent with other work showing that large disturbances can have cascading effects that can lead to continued habitat degradation. Effects of large disturbances are well known to have long-lasting impacts on coral communities (Done 1992, Connell et al. 1997, Connell 1997). Hydrodynamic energy from hurricanes is related to declines in reef structural complexity (Alvarez-Filip et al. 2009), low coral recruitment (Crabbe et al. 2008), and coral loss in the Caribbean over the last 30 yr (Gardner et al. 2005). On Caribbean reefs, the combined effects of disturbances, adult mortality, and recruitment limitations of framework-building scleractinian coral species have contributed to community shifts to small, weedy coral species, octocorals, or algae, thereby reducing potential for future structural complexity. This limitation is likely to have impacts beyond benthic communities, as complexity provided by both corals and underlying geologic structure contributes to system-wide biodiversity (Graham & Nash 2013), including fish communities (Graham et al. 2006), fisheries productivity (Graham 2014, Rogers et al. 2014), and mitigation of nearshore wave energy through attenuation by reefs (Ferrario et al. 2014).

After an extreme initial disturbance to the biological and structural complexity of a coral reef, projected recovery may be limited by ecological as well as hydrodynamic forcing. Our findings emphasize the importance of physical and biological limitations on juvenile coral survival and negative effects of multiple disturbances on community recovery. It is clear that without substrate stabilization, subsequent chronic hydrodynamic mobilization of unstable substrate can lead to prolonged or permanent habitat loss.

Acknowledgements. We thank Ryan Neve, Tom Moore, Greg Piniak, Brian Degan, Christine Buckel, Michael Nemeth, Katie Flynn, Sean Meehan, David Hurley, Pedro Rodriguez, and SeaVentures for their generous assistance. Constructive comments from 2 anonymous reviewers, Curt Storlazzi, Greg Piniak, Margaret Miller, Carolyn Currin, and Tom Moore improved the manuscript. This study was supported by NOAA National Centers for Coastal Ocean Science, NOAA Restoration Center, PADI Foundation, Sigma Xi Grant-In-Aid of Research, Duke University Graduate School, and the National Science Foundation (OCE-1435133). The scientific results and conclusions, as well as any views or opinions expressed herein, are those of the authors and do not necessarily reflect the views of NOAA or the Department of Commerce.

LITERATURE CITED

The following supplement accompanies the article

Understanding differential patterns in coral reef recovery: chronic hydrodynamic disturbance as a limiting mechanism for coral colonization

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Marine Ecology Progress Series 605: 135–150

1. Benthic cover

Benthic cover was compared between the rubble, pavement, restoration and reference sites with planar photographs of permanent quadrats. Benthic cover was quantified using manual classification in CoralNet software (Beijbom et al. 2012). Classification categories included crustose coralline algae (CCA), coral, cyanobacteria, Millepora spp. hydrocoral, macroalgae, octocoral, other, bare substrate, sponges, and turf algae. Percentage values per quadrat were adjusted to exclude points landing on survey hardware. 105 photos were analyzed; however, photographs from 2009 and 2011 were not included due to lack of representation at all survey sites.

Throughout the study, rubble dominated the cover of substrate at the rubble site, and consolidated substrate dominated at all other sites (Fig. S1A). Turf algae dominated biological cover at all sites and years, although some annual variability was evident (reference site: 56.9 - 67.4%; restoration site: 57.4 - 81.6%; pavement site: 62.1% - 99%; 83.9 - 91.5%; Fig. S1B). The reference site had considerably higher macroalgal cover (0.5 - 26.9%; restoration site: 0.95 - 8.6%; pavement: 0.3 - 2.2%; rubble site: 0 - 1.65%; Fig. S1B). Cyanobacterial mats were present at the pavement and restoration sites in 2012 and 2013, and highest in 2012 at the pavement site (6.2%; Fig. S1B). Octocoral cover was highest 2.7 - 19.9% at the reference site, followed by the
restoration site (5.7 - 16.5%), pavement site (0.2 - 6.1%), and rubble site (0 - 2.9%), slightly higher than scleractinian coral cover in the reference (4.8 - 10.3%), restoration (3.3 - 4.3%), pavement (0 - 1.6%), and rubble sites (0 - 0.2%; Fig. S1B), although cover measurement of octocorals is imprecise with a planar approach using photographs.

**Fig. S1A.** Mean cover (%) and standard error of substrate type for each sample level. No bar indicates zero cover. Years 2009 and 2011 are not included due to lack of field survey photographs at all levels.
**Fig. S1B.** Mean cover (%) and standard error of benthic cover for each sample level. No bar indicates zero cover. Years 2009 and 2011 are not included due to lack of field survey photographs at all levels.
2. Temperature comparison

Near-bottom water temperatures (0.5 m above substrate) were measured at rubble and pavement sites during the intensive oceanographic measurement period (13 June – 17 July 2014) using high precision thermistors (Seabird SBE-56). Near bottom mean temperatures at the pavement site were 28.80 ± 0.22 °C and at the rubble site were 28.66 ± 0.22 °C, although the distribution of pavement site temperatures was skewed toward higher temperatures (Fig S2).

![Graph showing temperature distribution](image)

**Fig. S2.** Distribution of near-bottom water temperatures from thermistors (SBE56, Seabird Electronics) deployed at the rubble and pavement sites.

3. Wave model calibration and comparisons with field data

A systematic set of wave model simulations was conducted to test the sensitivity of model output to the model parameters for bottom roughness \((k_w; \text{value range tested: 0.01 - 1})\), as used in Madsen et al. (1988), and the wave breaking coefficient \((\gamma_w; \text{value range tested: 0.5 - 1})\) based on the bore model of Battjes and Janssen (1978). We also tested wave model sensitivity to inclusion of tidal forcing in the open boundary conditions, and inclusion of whitecapping effects based on Komen et al. (1984). Based on comparison with the field measurements, the optimal
wave model parameters were found to be $k_w = 0.2$ and $\gamma_w = 0.6875$. Model output at the three study sites was sensitive to the choice of bottom roughness, but was relatively insensitive to $\gamma_w$, not unsurprising given the relatively deep water depths at those sites, but also suggesting that breaking in other parts of model domain did not have a large effect on wave conditions at the three study sites. Application of time-varying water level increased overall model skill at each site by a negligible amount; however, tidal variation was not included in the final model hindcast runs because these forcing data were not available during the 4-year hindcast period.

4. Derivation of equations for rubble motion from external forces

We modeled the potential mobility of unconsolidated rubble within the disturbed areas using a mechanistic approach to estimate hydrodynamic forces required to initiate rubble motion. This approach has been well developed in the literature to hindcast hydrodynamic forcing required to mobilize boulders during large wave events (Nandasena et al. 2011). Here, a clast is approximated by an idealized cuboid (defined by $a$, $b$, $c$-axes) that is acted upon by gravity ($F_g$), normal ($F_N$), lift ($F_L$), friction ($F_F$), drag ($F_D$), and inertial ($F_I$) forces, where $F_F = \mu F_N$ and $\mu$ is the coefficient of static friction (Fig. 3). Hydrodynamic forcing exceeding a critical threshold can initiate rubble clast motion by sliding or over-turning. Sliding occurs when

$$F_D(t) + F_I(t) > \mu [F_g - F_L(t)]$$

At the initiation of sliding, the forces on the left and righthand side of Eq. (1) are equal. Forces that are functions of velocity ($F_D$, $F_L$, and $F_I$) vary with time ($t$) over a wave period. Grouping together terms that are quadratic in velocity gives

$$[F_D(t) + \mu F_L(t)] + F_I(t) - \mu F_g = 0$$

The quadratic and lift forces due to waves are in phase with each other; however, they are both 90° out of phase with the inertial force. To estimate when rubble are mobilized, we are interested in when the total time-varying forces are maximum ($F_{TM}$). Following the derivation of Dean and Dalrymple (1991), but extending it to include the lift force (2) can be represented as

$$F_{D,m} + \mu F_{L,m} + \frac{F_{L,m}^2}{4(F_{D,m} + \mu F_{L,m})} - \mu F_g = 0$$  \hspace{1cm} (3)

where the maximum drag force is,

$$F_{D,m} = \frac{1}{2} \rho_w C_D A_D U_m^2$$  \hspace{1cm} (4)

the maximum lift force is,

$$F_{L,m} = \frac{1}{2} \rho_w C_L A_L U_m^2$$  \hspace{1cm} (5)

the maximum inertial force is,

$$F_{I,m} = \rho_w C_M V(U_m \omega)$$  \hspace{1cm} (6)

where $\omega = 2\pi / T$ is the wave frequency, $T$ is the wave period, and $V$ is the rubble volume. $\mu$, $C_D$, $C_L$, and $C_M$ are dimensionless coefficients of static friction, drag, lift, and mass respectively. $A_D$ is the cross-sectional area of a rubble clast perpendicular to the flow and $A_L$ is the area of a rubble clast parallel to the flow. The buoyant weight is,

$$F_g = \Delta \rho V g$$  \hspace{1cm} (7)

where $\Delta \rho = \rho_r - \rho_w$, $\rho_w$ is the density of seawater (1023 kg m$^{-3}$), $\rho_r$ is rubble density, and $g$ is acceleration due to gravity.

The water velocity can be decomposed into contributions from current and waves (Grant & Madsen 1979)
Here we consider the case where currents are assumed to be constant (since current data are limited to the duration of the short-term field deployment), but waves vary with time (since wave hindcast data span 4-years),

\[ U_w(t) = U_c(t) \]  

(9)

Substituting the above expressions for the forces into (3), and noting that the currents do not contribute to the inertial force

\[
\begin{align*}
\frac{1}{2} \rho \omega^2 DC_A (U_c + U_w)^2 + \mu \frac{1}{2} \rho \omega^2 DC_A (U_c + U_w)^2 \\
+ \left( \frac{\rho \omega^2 DC_A}{4} \right) (U_c + U_w)^2 + \mu \frac{1}{2} \rho \omega^2 DC_A (U_c + U_w)^2 \\
- \mu \Delta \rho V g = 0
\end{align*}
\]

(10)

We solve for \( U_w \), the threshold wave orbital velocity for rubble sliding; all other variables have known values. To simplify notation, let

\[
a = (\rho \omega^2 DC_A + \mu \rho \omega^2 DC_A) \\
b = \frac{\rho \omega^2 DC_A}{\rho \omega^2 DC_A + \mu \rho \omega^2 DC_A} \\
c = -2 \mu \Delta \rho V g \\
d = U_c
\]

(11) \quad (12) \quad (13) \quad (14)

So that

\[ a(d + U_w)^4 + b(U_w)^2 + c(d + U_w)^2 = 0 \]  

(15)

Expanding out polynomials and collecting like order terms yields a \( 4^{th} \) order polynomial equation

\[ (a)U_w^4 + (4ad)U_w^3 + (6ad^2 + b + c)U_w^2 + (4ad^3 + 2cd)U_w + (ad^4 + cd^2) = 0 \]  

(16)

that we solved numerically and retained only physically realistic roots \( i.e., \) positive real numbers.)
As an aside, note that for the case of no background mean current \( (U_c = 0) \), equation (16) reduces to

\[
U_m = \sqrt{\frac{\mu \Delta \rho V g - \left( \frac{C_M V \omega}{2(C_D A_D + \mu C_L A_L)} \right) C_M V \omega}{\frac{1}{2} (C_D A_D + \mu C_L A_L)}}
\]

Equation (17) is similar to Buckley et al. (2012)'s equation 6 for the threshold velocity of rubble motion, but with the addition of the term in grey, which represents a coefficient that accounts for the phasing required for the combined drag, lift, and inertial forces to be maximum. Simplifying (17) slightly gives an expression representing the threshold water velocity for rubble sliding,

\[
u_{\text{slide,m}} = \sqrt{\frac{2\mu \Delta \rho V g - \left( \frac{C_M V \omega}{2(C_D A_D + \mu C_L A_L)} \right)^2}{(C_D A_D + \mu C_L A_L)}}
\]

If \( F_I > 2(F_D + F_L) \), then the total force is dominated by the inertial force, and (3) is replaced with the following equations (Dean and Dalrymple 1991)

\[
F_{I,m} = \mu F_g
\]

\[
C_M V U_m \omega = \mu \Delta \rho V g
\]

\[
U_m = \frac{\mu \Delta \rho g}{C_M \omega}
\]

And so for cases where \( F_I > 2(F_D + F_L) \)

\[
u_{\text{slide}} = \frac{\mu \rho g}{C_M \omega}
\]
Next we examine the case of rubble motion due to overturning. From the free-body diagram (Fig 3), rubble overturning occurs when the sum of the overturning moments exceed the restoring moment

$$F_D l_D + F_l l_I + F_L l_L > F_g l_g$$  \hspace{1cm} (23)

where the moment arm for the drag and inertial forces are assumed to be the same (half the rubble height),

$$I_D = I_I = I_n / 2 \equiv I_H$$ \hspace{1cm} (24)

and the moment arm for the lift and gravitational forces are assumed to be the same (half the rubble width)

$$I_L = l_g = I_w / 2 \equiv I_W$$ \hspace{1cm} (25)

Substituting (24) and (25) into (23) and grouping terms

$$(F_D + F_I) I_H + (F_L - F_g) I_W = 0$$ \hspace{1cm} (26)

As in the derivation above for sliding, we are interested in time-varying overturning moments and so the internal force must be modified to account for relative phasing with drag and lift forces. Following Dean and Dalrymple (1991)

$$F_i = \frac{F_{I,m}^2}{4(F_{D,m} + F_{L,m})}$$ \hspace{1cm} (27)

Substituting (27) and expressions for drag, lift, inertial and gravitational forces into (26) and simplifying gives

$$\left(\rho_w^2 C_D^2 A_D^2 l_H + \rho_n^2 C_D C_L A_D A_L l_H + \rho_n^2 C_L^2 C_D A_D A_L l_H + \rho_n^2 C_L^2 A_L^2 I_W \right) U_m^4 + \left(\rho_n^2 C_D^2 V^2 \alpha^2 l_H - 2 \rho_n C_D A_D \Delta \rho V g l_w - 2 \rho_n C_L A_L \Delta \rho V g l_w \right) U_m^2 = 0$$ \hspace{1cm} (28)

Then the overturning equation has the general quartic form where $\alpha$ and $\beta$ are non-integers.
which was solved numerically for $U_m$, and only physically realistic roots retained.

\[
\alpha U_m^4 + \beta U_m^2 = 0
\]  

(29)

5. Quantification of rubble morphology and sensitivity analyses

Rubble morphology was quantified to make force estimates in the rubble disturbance area by measuring surface area and buoyant weight from a random selection of rubble clasts collected from the rubble site in 2013 and 2014. Projected areas of individual clasts ($n = 102$) were measured along three dimensions: $a$-axis (top, length), $b$-axis (side, width) and $c$-axis (end, height) using orthogonal digital images and image analysis software (ImageJ; Rasband 1997; Fig S3). Area normal to flow ($A_{p}$) was considered to be along the $a$-axis; area perpendicular to flow ($A_{t}$) was considered to be along the plane of the $b$- and $c$-axes. Rubble volume ($V$) was calculated as the product of $a$, $b$, and $c$-axis lengths, where due to the irregular rubble shapes, each axis length was approximated as the square root of the corresponding area ($a-c$). Weights for dry and submerged rubble were measured using a digital force sensor (Vernier, Beaverton, OR), and buoyant weight was then calculated following Jokiel et al. (1978). In the general case, we use the median values of the range of measurements to determine effects of rubble properties on threshold bottom velocities for overturning ($u_{over}$) and sliding ($u_{slide}$) using (18) and (29). Sensitivity analyses were conducted with rubble property quantiles (5, 10, 25, 50, 75, 95%) from the distribution of each measurement (Fig S4). The smallest size class, at 5% rubble quartile, typically produced NA mobilization values.
FIG. S3. Measured rubble A) area, B) volume, and C) density. In A) rubble area perpendicular to flow ($A_{D}$) is dark grey, and rubble area normal to the flow ($A_{L}$) is light grey.
In an expanse of reef rubble, proximity of multiple clasts likely impacts hydrodynamic forces affecting any given clast, due to enhanced friction from interlocking or by reduced area exposed to flow due to sheltering from adjacent objects. We addressed this by including a blocking term that accounts for reduction in $A_D$, following Storlazzi et al (2005). This term has a value from 0-1, corresponding to flow totally blocked to no blockage. We applied a value of 0.5 with recognition that a wide array of conditions is likely to exist in situ. Sensitivity analyses show little variability with a blockage value range from 0.1 to 0.9 (Fig. S4).

**Fig. S4.** Bottom orbital velocity ($u_b$) estimated to mobilize rubble through sliding or over turning, based on the distribution of rubble sizes (%) collected from within the rubble disturbance.

Summary statistics of rubble clast dimensions showed variability primarily in size and volume; sensitivity studies showed that the minimum bottom orbital velocity required to mobilize rubble was influenced by rubble size and exposure to flow (Fig. S5). Drag, lift, inertia and frictional forces estimated at each site were not significantly different between sites (Fig. S5), although the data suggest that the pavement sites had slightly higher forcing.
FIG. S5. Drag ($F_d$), lift ($F_L$), inertial ($F_I$) and frictional ($F_F$) forces estimated at the rubble and pavement sites, as calculated from mean rubble statistics and wave model output from 2010 - 2014. The box plot components are: solid line, median; box, interquartile range (IR); whiskers, most extreme point less than 1.5 times the IR from the median.
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