

# Dynamic Models for Yielding and Friction Hysteresis

## CEE 541. Structural Dynamics

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In materials or elements with hysteresis, the response to a cycle reciprocating forcing depends on the forcing history — for *any* reciprocating forcing of a sufficiently large amplitude [10]. Hysteretic behavior is commonly depicted as loops in graphs (2D plots) of periodic (oscillatory) output vs. periodic input. In rate-dependent hysteresis, the size and shape of the hysteresis loop changes with the rate or frequency of the input. If the loop collapses to a function (e.g., a curved line) for any input (e.g., quasi-static), then the system is *not* hysteretic [10]. Hysteresis implies a non-linear relationship between inputs and outputs: differential equation models for hysteresis must be nonlinear; convolution models for hysteresis must be nonhomogeneous. Linear visco-elastic materials are rate-dependent but are *not hysteretic* because forces and displacements are proportional in the limit of quasi-static loading. This document describes dynamic hysteresis models that are members of the Duhem [9, 15, 17, 20] class of nonlinear ordinary differential equations

$$\dot{z}(t) = f(z(t), u(t)) g(\dot{u}(t)) , \quad z(0) = z_o . \quad (1)$$

Duhem models can be used to model the kinds of rate-independent hysteresis representative of material yielding and stick-slip friction.

Consider the following nonlinear ordinary differential equation

$$\dot{z}(t) = \dot{u}(t) - |\dot{u}(t)| z^\eta(t) , \quad z(0) = z_o . \quad (2)$$

The variable  $z(t)$  represents a level of force, normalized to be in the range  $-1 < z(t) < 1$ , so you can think of  $z$  as a force level normalized by a yield force, or friction sliding force. The variable  $u(t)$  represents a level of displacement, normalized by a yield displacement, so you can think of  $u$  as a ductility ratio. In this way, equation (2) relates the rate of change of force to the rate of change of displacement (i.e., velocity). The exponent  $\eta$  is restricted to be odd, for the time being.

Noting that  $|\dot{u}| = \dot{u} \operatorname{sgn}(\dot{u})$ , equation (2) may be re-written as,

$$\dot{z} = (1 - z^\eta \operatorname{sgn}(\dot{u})) \dot{u} , \quad z(0) = z_o \quad (3)$$

and the slope of the force-displacement relationship is

$$\frac{dz}{du} = \frac{\dot{z}}{\dot{u}} = 1 - z^\eta \operatorname{sgn}(\dot{u}) \quad (4)$$

From this expression it is easy to see that:

- When the force is zero ( $z = 0$ ), the dimensionless stiffness is 1 (the dimensional stiffness is the yield force divided by the yield displacement).
- As the force approaches the yield force ( $z \rightarrow 1, \dot{u} > 0$  or  $z \rightarrow -1, \dot{u} < 0$ )  $dz/du$  approaches zero.
- When the velocity is positive,  $dz/du = 1 - z^\eta$ , and when the velocity is negative,  $dz/du = 1 + z^\eta$ .
- $dz/du \geq 0$  ;  $\text{sgn}(\dot{z}) = \text{sgn}(\dot{u})$  ; and  $(\dot{z})(\dot{u}) \geq 0$  .

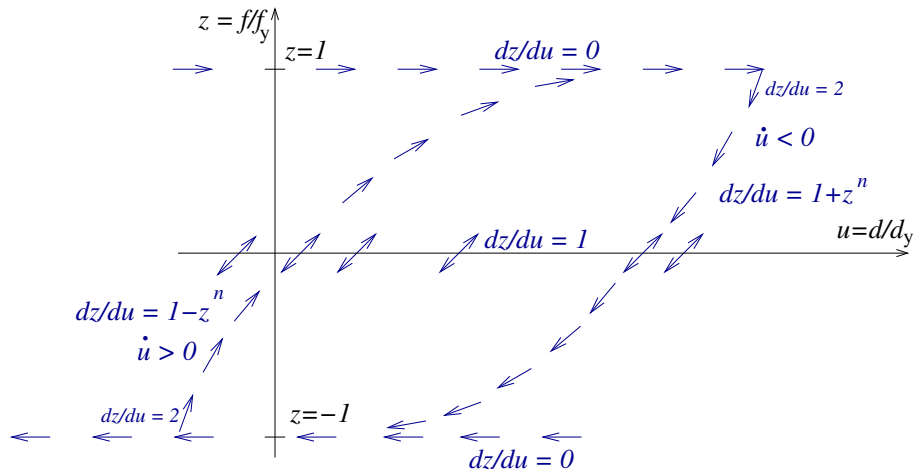


Figure 1. The vector field of  $dz/du$  for  $\dot{z} = (1 - z^\eta \text{sgn} \dot{u}) \dot{u}$  depends on  $\text{sgn} \dot{u}$ . Larger values of  $\eta$  result in a sharper “knee.”

### Extension 1

Replacing  $z^\eta$  with  $|z|^\eta \text{sgn}(z)$  in equation (3), and noting that  $\text{sgn}(a)\text{sgn}(b) = \text{sgn}(ab)$ ,

$$\dot{z} = (1 - |z|^\eta \text{sgn}(\dot{u}z)) \dot{u} , \quad (5)$$

which allows the exponent  $\eta$  to be any positive value.

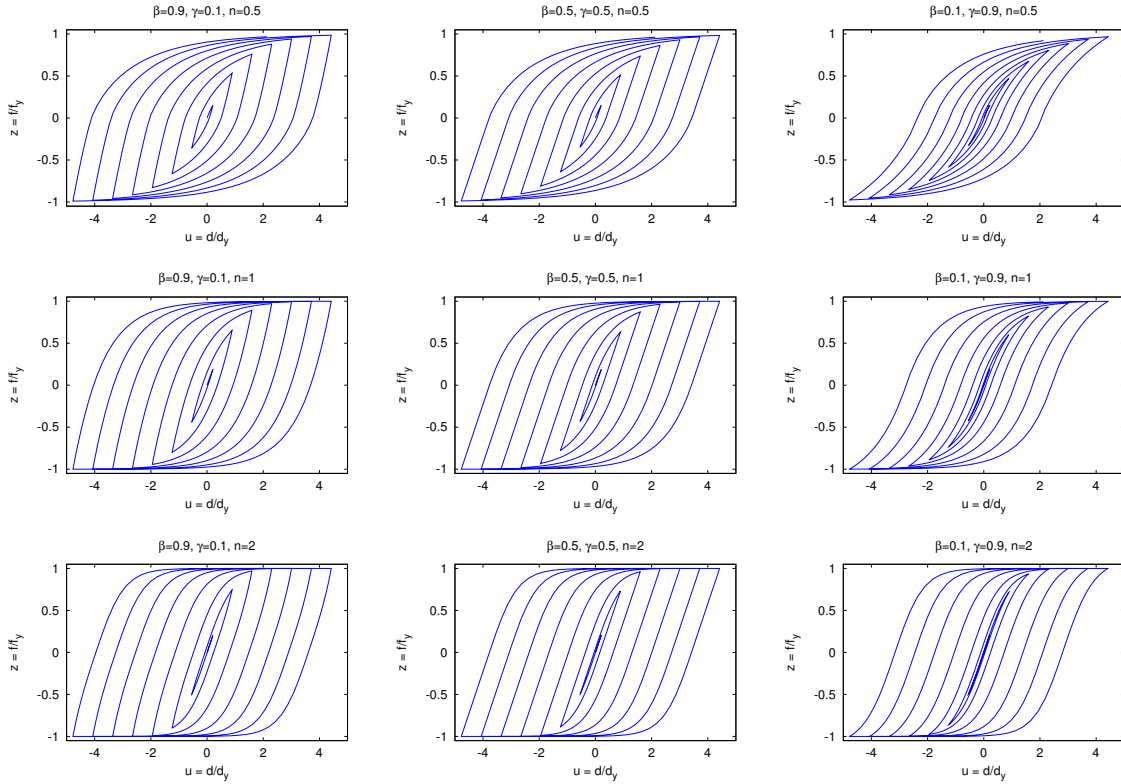
### Extension 2

Adding parameters  $\beta$  and  $\gamma$  in equation (5), as follows,

$$\dot{z} = (1 - |z|^\eta (\beta \text{sgn}(\dot{u}z) + \gamma)) \dot{u} , \quad (6)$$

allows for a wide range of hysteretic forms, as shown in figure 2.

If  $\beta + \gamma = 1$ , then  $-1 < z < 1$ . If  $\eta > 0$ ,  $\gamma > 0$  and  $-\gamma < \beta < \gamma$  the model respects the Second Law of Thermodynamics [1, 11, 14].


 Figure 2. Dependence of hysteretic shape on  $\beta$ ,  $\gamma$ , and  $\eta$ .

### Extension 3

Adding a parameter  $A$ ,

$$\dot{z} = (A - |z|^\eta (\beta \operatorname{sgn}(\dot{u}z) + \gamma)) \dot{u}, \quad (7)$$

allows for scaling. If  $\beta + \gamma = A$  then  $-A < z < A$ . This is the ‘‘Bouc-Wen’’ model for hysteresis [5, 6, 14, 26] and is an element of the Duhem class of hysteresis models [17, 25].

### Extension 4

Isotropic bi-axial hysteretic behavior may be modeled in orthogonal directions  $x$  and  $y$  by coupling the hysteretic variables and velocities [12, 23].

$$\dot{z} = a(z, \dot{u}) z + b \dot{u} \quad (8)$$

where  $z = [z_x, z_y]$ ,  $\dot{u} = [\dot{u}_x, \dot{u}_y]$ , and

$$a(z, \dot{u}) = -[\beta(|z_x \dot{u}_x| + |z_y \dot{u}_y|) + \gamma(z_x \dot{u}_x + z_y \dot{u}_y)] (z_x^2 + z_y^2)^{(\eta-2)/2} \quad (9)$$

### Extension 5

Hysteretic behavior with a non-zero post yield stiffness may be simulated by combining equation (7) or (8) with

$$f(t) = f_y((1 - \kappa)z(t) + \kappa u(t)), \quad (10)$$

where  $f_y$  is a yield force level,  $u$  is the displacement divided by the yield displacement (the ductility), and  $\kappa$  is the ratio of the post yield stiffness to the pre-yield stiffness. When using (10), set  $A = 1$  and  $\beta + \gamma = 1$ .

## Stick-Slip Friction

The Dahl friction model [8],

$$\dot{z} = (1 - z \operatorname{sgn}(\dot{u}))^\eta \dot{u} \quad (11)$$

is equivalent to the Bouc model for  $\eta = 1$  (which is the usual choice for  $\eta$  in friction modeling). The “LuGre” friction model [7] is an extension of the Dahl model which captures the Stribeck (“stick-slip”) effect [2]. The LuGre model is an element of the Duhem class of hysteresis models [17, 25].

$$\dot{z} = (1 - z \operatorname{sgn}(\dot{u})/g(\dot{u})) \dot{u} \quad (12)$$

$$k_o g(\dot{u}) = F_C + (F_S - F_C) e^{-(\dot{u}/v_S)^2} \quad (13)$$

$$f(t) = k_o z(t) + c_1 \dot{z}(t) + c_2 \dot{u}(t) \quad (14)$$

In the LuGre model the velocity  $\dot{u}$  is not normalized by a pre-slip displacement. For high values of pre-slip stiffness,  $k_o$ , the LuGre model can require very small time steps for numerical stability.

Table 1. Representative parameter values for the LuGre friction model [7].

Parameter	Definition	Value	Unit
$k_o$	pre-slip stiffness	$10^4$	N/m
$c_1$	friction rate effect	$\sqrt{10^4}$	Ns/m
$c_2$	viscous rate effect	0.4	Ns/m
$F_C$	Coulomb sliding friction force	1	N
$F_S$	Stribeck sticking friction force	1.5	N
$v_S$	Stribeck velocity	0.001	m/s

## Degrading Behavior

To model the accumulation of damage [4], strength  $f_y$  and other model parameters may be linked to a damage accumulation index  $\mathcal{D}$ , where  $\dot{\mathcal{D}} \approx |\dot{u} - \dot{z}|$  and

$$f_y(t) = \frac{f_y(0)}{1 + \mathcal{D}(t)/a_{(f_y)}}, \quad (15)$$

where  $a_{(f_y)}$  is a positive constant.

## Further Extensions

Further generalizations to these equations can account for orthotropic behavior, degrading stiffness, and pinching hysteresis, as described in the references below.

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